

Analysis on infinite configuration spaces, and L^2 -Betti numbers associated with infinite particle systems

Alexei Daletskii

University of York, UK

Based on joint works with
S Alberverio, A. Kalyuzhnyi, Yu. Kondratiev,
E. Lytvynov, Yu. Samoilenko

Introduction: What is a Configuration Space?

X - topological space (e.g. \mathbb{R}^d or Riemannian manifold)

Γ_X - space of locally finite subsets (configurations) in X :

$$\Gamma_X = \{ \gamma \subset X : |\gamma \cap \Lambda| < \infty, \Lambda \text{ compact} \}$$

Introduction: What is a Configuration Space?

X - topological space (e.g. \mathbb{R}^d or Riemannian manifold)

Γ_X - space of locally finite subsets (configurations) in X :

$$\Gamma_X = \{ \gamma \subset X : |\gamma \cap \Lambda| < \infty, \Lambda \text{ compact} \}$$

- Statistical mechanics (models of gases) - Dobrushin 56, 70; Ruelle 63, 70)
- Quantum Field Theory - Goldin, Grodnik, Powers, Sharp 75
- Representation Theory - Gelfand, Graev, Vershik 75
- Probability (theory of point processes) - Föllmer 75; Preston 76, 79; Georgii 76
- Topology - Fadell 62; Bödingheimer (spaces of *finite configurations*)

Introduction: What is a Configuration Space?

X - topological space (e.g. \mathbb{R}^d or Riemannian manifold)

Γ_X - space of locally finite subsets (configurations) in X :

$$\Gamma_X = \{ \gamma \subset X : |\gamma \cap \Lambda| < \infty, \Lambda \text{ compact} \}$$

- Statistical mechanics (models of gases) - Dobrushin 56, 70; Ruelle 63, 70)
- Quantum Field Theory - Goldin, Grodnik, Powers, Sharp 75
- Representation Theory - Gelfand, Graev, Vershik 75
- Probability (theory of point processes) - Föllmer 75; Preston 76, 79; Georgii 76
- Topology - Fadell 62; Bödingheimer (spaces of *finite configurations*)

Γ_X as *infinite dimensional manifold* - Albeverio, Kondratiev, Röckner 96

Topology and measures

$$\Gamma_X \ni \gamma \simeq \sum_{x \in \gamma} \delta_x \in \mathcal{M}(X), \quad \delta_x - \text{Dirac measure}$$

Vague topology on Γ_X - weak topology induced from $\mathcal{M}(X)$

$$\Gamma_X \ni \gamma \simeq \sum_{x \in \gamma} \delta_x \in \mathcal{M}(X), \quad \delta_x \text{- Dirac measure}$$

Vague topology on Γ_X - weak topology induced from $\mathcal{M}(X)$

Poisson Measure π_σ (with reference measure σ)

- 1 Consider space of finite configurations in a compact $\Lambda \subset X$

$$\Gamma_\Lambda = \cup_n \widetilde{\Lambda}^n / S_n;$$

define *probability measure* π_σ^Λ on Γ_Λ via formula

$$\pi_\sigma^\Lambda = e^{-\sigma(\Lambda)} \sum_n \frac{1}{n!} \sigma^n \quad (\sigma \text{- infinite measure on } X)$$

- 2 Define

$$\pi_\sigma = \lim \pi_\sigma^\Lambda, \quad \Lambda \uparrow X.$$

Quasi-invariance and gradient

- $Diff_0(X)$ - group of compactly supported diffeomorphisms of X - acts on Γ_X :

$$\{\dots, x, y, z, \dots\} \mapsto \{\dots, \varphi(x), \varphi(y), \varphi(z), \dots\}, \quad \varphi \in Diff_0(X);$$

Main fact: π_σ is $Diff_0(X)$ - quasi-invariant.

- Γ -gradient: for $v \in Vect_0(X)$, $F : \Gamma_X \rightarrow \mathbb{R}$ set

$$\begin{aligned} \nabla_v^\Gamma F(\gamma) &: = \frac{d}{dt} F(\varphi_t^v \gamma)_{t=0} \\ &= \sum_{x \in \gamma} (\nabla_x F(\gamma), v(x))_{T_x X} \end{aligned}$$

and

$$\nabla^\Gamma F(\gamma) := (\nabla_x F(\gamma))_{x \in \gamma} \in \bigoplus_{x \in \gamma} T_x X$$

- Tangent space (Gelfand, Graev, Vershik 75):

$$T_\gamma \Gamma_X := \bigoplus_{x \in \gamma} T_x X$$

Integration by parts formula

Albeverio, Kondratiev, Röckner 96

$$\int_{\Gamma_X} \nabla_v^\Gamma F(\gamma) \pi(d\gamma) = - \int_{\Gamma_X} F(\gamma) B_\pi^v(\gamma) \pi(d\gamma),$$
$$B_\pi^v(\gamma) = \langle \beta_\sigma^v, \gamma \rangle$$

for local function F (that is $F(\gamma) = F(\gamma \cap \Lambda)$)

$B_\pi^v(\gamma)$ - logarithmic derivative of π along v ,

β_σ^v - logarithmic derivative of σ along v ($= \operatorname{div} v$ if σ is Riemannian volume)

General notation: $\langle f, \gamma \rangle := \sum_{x \in \gamma} f(x)$, $f \in C_0(X)$

Laplace (Dirichlet) operator

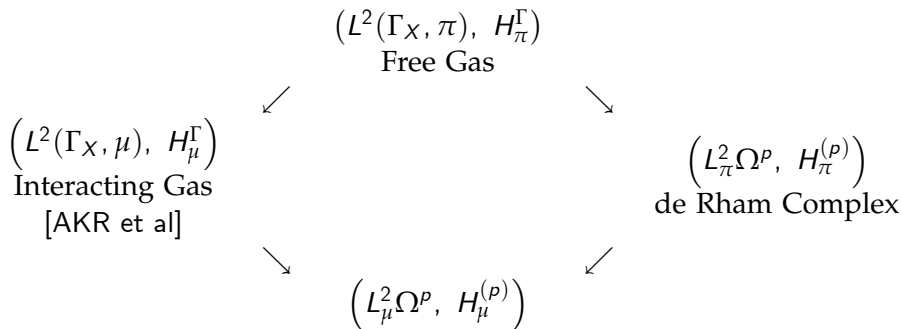
Dirichlet form: for local functions F and G on Γ_X

$$\begin{aligned}\mathcal{E}_\pi(F, G) &: = \int_{\Gamma_X} \left(\nabla^\Gamma F(\gamma), \nabla^\Gamma G(\gamma) \right)_{T_\gamma \Gamma_X} \pi(d\gamma) \\ &\stackrel{IBP}{=} - \int_{\Gamma_X} \left(H_\pi^\Gamma F \right) (\gamma) G(\gamma) \pi(d\gamma)\end{aligned}$$

Dirichlet operator $H_\pi^\Gamma F(\gamma) = \sum_{x \in \gamma} (\Delta_x + \beta_\sigma \nabla_x) F(\gamma)$ - self-adjoint operator in $L^2(\Gamma_X, \pi)$

- H_π^Γ is essentially self-adjoint on smooth local functions
- $\text{Ker} H_\pi^\Gamma = \{const\}$
- $L^2(\Gamma_X, \pi) \simeq \text{Fock}(L^2(X, \sigma))$ and $H_\pi^\Gamma \simeq$ second quantization of $\Delta + \beta_\sigma \nabla$
- $\exp(-tH_\pi^\Gamma)$ - Markov semigroup, corresponding stochastic process is *independent Brownian motion of particles*

Where to go now?



μ is Gibbs or cluster Gibbs measure

Differential forms on configuration space

Reminder: $T_\gamma \Gamma_X = \bigoplus_{x \in \gamma} T_x X$

- 1-form (vector field) $\Gamma_X \ni \gamma \mapsto \omega(\gamma) \in T_\gamma \Gamma_X$
- p -form

$$\Gamma_X \ni \gamma \mapsto \omega(\gamma) \in \Lambda^p T_\gamma \Gamma_X$$

- We work with square-integrable forms $\omega \in L^2_\pi \Omega^p$.
Define Hodge differential

$$d_p : L^2_\pi \Omega^p \rightarrow L^2_\pi \Omega^{p+1}$$

by $d_p \omega(\gamma) = \text{anti-symmetrization } (\nabla^\Gamma \omega(\gamma))$.

$$d_p^* : L^2_\pi \Omega^{p+1} \rightarrow L^2_\pi \Omega^p$$

is densely defined (because of IBP)

- Hodge-de Rham Laplacian

$$\mathbb{H}^{(p)} := d_p d_p^* + d_{p-1}^* d_{p-1},$$

self-adjoint in $L^2_\pi \Omega^p$

Decomposition of Hodge-de Rham Laplacian

Albeverio, AD, Lytvynov:

1 \exists isometry $\mathcal{I}^p : L^2_\pi \Omega^p \rightarrow L^2(\Gamma_X, \pi) \otimes [\bigoplus_{m=1}^p L^2_\sigma \Omega^p(X^m)]$

2

$$\mathbf{H}_\pi^{(p)} \stackrel{\mathcal{I}^p}{\simeq} H_\pi^\Gamma \otimes \mathbf{1} + \mathbf{1} \otimes \left[\sum_{m=1}^p H_{X^m}^{(p)} \right],$$

where $H_{X^m}^{(p)}$ is Hodge-de Rham on X^m

3 Künneth formula

$$\text{Ker } \mathbf{H}_\pi^{(p)} \stackrel{\mathcal{I}^p}{\simeq} \bigoplus_{s_1, \dots, s_d} \left(\mathcal{H}_X^{(1)} \right)^{\diamond_{s_1}^1} \otimes \dots \otimes \left(\mathcal{H}_X^{(d)} \right)^{\diamond_{s_d}^d},$$

$\mathcal{H}_X^{(m)} := \text{Ker } H_X^{(m)}$, $d := \dim X$,

$$\diamond_s = \begin{cases} \widehat{\otimes}, & m \text{ even} \\ \wedge, & m \text{ odd} \end{cases}$$

Corollary: let $\beta_m := \dim \mathcal{H}_X^{(m)} < \infty$. Then

$$\dim \text{Ker } \mathbf{H}_\pi^{(p)} = \sum_{s_1, \dots, s_d} \beta_1^{(s_1)} \dots \beta_d^{(s_d)},$$

$$\beta_m^{(s)} = \binom{\beta_m}{s}, \quad m \text{ odd, and } \beta_m^{(s)} = \binom{\beta_m + s - 1}{s}, \quad m \text{ even.}$$

Example: manifold with cylinder end $X = M \cup (N \times \mathbb{R}^1)$, M compact with boundary N .

If $d = 2$ then $\beta_0 = \beta_2 = 0$, $\beta_1 < \infty$. Then

$$\dim \text{Ker } \mathbf{H}_\pi^{(p)} = \begin{cases} \binom{\beta_1}{p} & \text{if } p \leq \beta_1 \\ 0 & \text{if } p > \beta_1 \end{cases}$$

Remark: In general, $\dim \mathcal{H}_X^{(m)} = \infty$ because X is not compact

L2-Betti numbers (M. Atiyah 76)

Assumption: $\exists G \subset \text{Iso } X$, infinite discrete; $M = X/G$ - compact manifold. **Example:** $X = \mathbb{H}^d$ -hyperbolic space

$$\text{group action} \quad G \in g \mapsto T_g \in \mathcal{B}(L^2\Omega^p(X)).$$

Commutant of G -action

$$\mathcal{A} := \{T_g, g \in G\}' \subset \mathcal{B}(L^2\Omega^p(X)) -$$

von Neumann algebra (II_∞ factor if G is "strongly" non-commutative).

Orthogonal projection

$$P : L^2\Omega^p(X) \rightarrow \mathcal{H}_X^{(p)}.$$

- $P \in \mathcal{A}$, L^2 -Betti number $b_n = \dim_G \mathcal{H}_X^{(p)} := \text{Tr}_{\mathcal{A}} P < \infty$
- $b_n = 0$ iff $\dim \mathcal{H}_X^{(m)} < \infty$
- $\text{Index}_{\mathcal{A}}(d + d^*) := \sum (-1)^p b_p = \chi(M)$ - Euler characteristic of M

Traces of projections in tensor products of factors

$\mathcal{H} \subset \mathfrak{X}$ - Hilbert spaces; $P : \mathfrak{X} \rightarrow \mathcal{H}$ - orthogonal projection;
assume $P \in \mathcal{A}$ - some von Neumann algebra $\neq \mathcal{B}(\mathcal{H})$, $\text{Tr}_{\mathcal{A}} P < \infty$.

Define

$$P_s^{(n)} : \mathfrak{X}^{\otimes n} \rightarrow \mathcal{H}^{\widehat{\otimes} n}, \quad P_a^{(n)} : \mathfrak{X}^{\otimes n} \rightarrow \mathcal{H}^{\wedge n}.$$

$P_s^{(n)}, P_a^{(n)} \in \{\mathcal{A}^{\otimes n}, \{U_\alpha, \alpha \in S_n\}\}'' := \mathcal{A}^{(n)}$ ($\neq \mathcal{A}^{\otimes n}$ in general).

Traces of projections in tensor products of factors

$\mathcal{H} \subset \mathfrak{X}$ - Hilbert spaces; $P : \mathfrak{X} \rightarrow \mathcal{H}$ - orthogonal projection;
assume $P \in \mathcal{A}$ - some von Neumann algebra $\neq \mathcal{B}(\mathcal{H})$, $\text{Tr}_{\mathcal{A}} P < \infty$.

Define

$$P_s^{(n)} : \mathfrak{X}^{\otimes n} \rightarrow \mathcal{H}^{\widehat{\otimes} n}, \quad P_a^{(n)} : \mathfrak{X}^{\otimes n} \rightarrow \mathcal{H}^{\wedge n}.$$

$$P_s^{(n)}, P_a^{(n)} \in \{\mathcal{A}^{\otimes n}, \{U_\alpha, \alpha \in S_n\}\}'' := \mathcal{A}^{(n)} \quad (\neq \mathcal{A}^{\otimes n} \text{ in general}).$$

Theorem

1) \mathcal{A} - II_1 -factor. Then

$$\begin{aligned} \mathcal{A}^{(n)} &= W^*(\mathcal{A}^{\otimes n}, S_n) \quad (\text{cross-product}); \\ \text{Tr}_{\mathcal{A}} P_s^{(n)} &= \text{Tr}_{\mathcal{A}} P_a^{(n)} = \frac{(\text{Tr}_{\mathcal{A}} P)^n}{n!}. \end{aligned}$$

2) \mathcal{A} - II_∞ -factor, $\mathcal{A} = \mathcal{B}(\mathbb{C}^m) \otimes \mathcal{M}$, \mathcal{M} - II_1 -factor. Then

$$\mathcal{A}^{(n)} = \mathcal{B}(\mathbb{C}^m)^{\otimes n} \otimes W^*(\mathcal{M}^{\otimes n}, S_n).$$

L2-Betti numbers of Poisson configuration spaces

$$b_p(\Gamma_X) := \dim_G \text{Ker } \mathbf{H}_\pi^{(p)} = \sum_{s_1, \dots, s_d} \frac{b_1^{s_1}}{s_1!} \cdots \frac{b_d^{s_d}}{s_d!}, \quad d = \dim X - 1$$

Example. $X = \mathbb{H}^{d+1}$ - hyperbolic space,

$$b_p(X) = \begin{cases} 0, & p \neq \frac{1}{2}(d+1) \\ c, & p = \frac{1}{2}(d+1) \end{cases}, \quad p = 1, 2, \dots, d.$$

Then

$$b_p(\Gamma_X) = \begin{cases} 0, & p \neq \frac{1}{2}(d+1)s \\ \frac{c^s}{s!}, & p = \frac{1}{2}(d+1)s \end{cases}, \quad s = 1, 2, \dots$$

Index of the Dirac operator

$$D := d + d^* : L_\pi^2 \Omega^{\text{even}} \rightarrow L_\pi^2 \Omega^{\text{odd}}$$

$$\begin{aligned} \text{ind}_G &= \dim_G \text{Ker } D - \dim_G \text{Ker } D^* \\ &= \sum_p (-1)^p b_p(\Gamma_X) = \exp \chi(X/G). \end{aligned}$$

Gibbs measures on configuration spaces

$v : \mathbb{R} \rightarrow \mathbb{R}$; ρ - distance in X ; pair potential $V(x, y) = v(\rho(x, y))$;

$$\text{Energy} \quad E(\gamma) = \sum_{x, y \in \gamma} V(x, y), \quad \gamma \in \Gamma_X$$

$$\text{Gibbs measure} \quad \mu(d\gamma) = \frac{1}{Z} \exp(-E(\gamma)) \pi(d\gamma).$$

- μ - probability measure on Γ_X - well-defined (*conditions on potential V*);
- in general, μ is not unique (*phase transitions*);
- μ is $\text{Diff}_0(X)$ -quasi-invariant;
- μ -symmetric Hodge-de Rham Laplacian in $L^2(\Gamma_X, \mu)$ has *complicated structure*.

Random Witten Laplacian on X

Define measure σ_γ ($\gamma \in \Gamma_X$) on X :

$$\sigma_\gamma(dx) = \exp(-E_\gamma(x)) dx, \text{ where } E_\gamma(x) = \sum_{y \in \gamma} V(x, y).$$

Witten Laplacian $H_{\sigma_\gamma}^{(p)}$ associated with σ_γ :

symmetrization of the classical Hodge-de Rham $H^{(p)}$ in $L^2_{\sigma_\gamma} \Omega^p(X)$.

$$\begin{aligned} L^2_{\sigma_\gamma} \Omega^p(X) &\sim L^2 \Omega^p(X) \\ H_{\sigma_\gamma}^{(p)} &\sim H_\gamma^{(p)} = H^{(p)} + W_\gamma^{(p)}, \end{aligned}$$

where

$$W_\gamma^{(p)}(x) : (T_x X)^{\wedge p} \rightarrow (T_x X)^{\wedge p},$$

$$W_\gamma^{(p)}(x) = \|\nabla E_\gamma(x)\|^2 \text{id} + \Delta E_\gamma(x) \text{id} + (\nabla^2 E_\gamma(x))^{\wedge p}.$$

$H_\gamma^{(p)}$ is self-adjoint positive operator in $L^2 \Omega^p(X)$

Trace of the semigroup

Goal: to define (von Neumann) trace of the semigroup $e^{-tH_\gamma^{(p)}}$

Framework: $G \subset \text{Iso } X$, infinite discrete; $M = X/G$ - compact manifold;

σ_γ is G -invariant

Action of G on Γ_X :

$$\gamma = \{\dots, x, y, z, \dots\} \mapsto g\gamma = \{\dots, gx, gy, gz, \dots\}, \quad g \in G;$$

T_g - associated (unitary) representation of G in $L^2\Omega^p(X)$;

$$T_g H_\gamma^{(p)} T_g^{-1} = H_{g\gamma}^{(p)}.$$

Thus $e^{-tH_\gamma^{(p)}} \notin \mathcal{A} := \{T_g, g \in G\}'$ for μ -a.a. $\gamma \in \Gamma_X$
(since $\mu(\{\gamma : g\gamma = \gamma \ \forall g \in G\}) = 0$).

Define

$$\mathcal{C} = \{A : \Gamma_X \rightarrow \mathcal{B}(L^2\Omega^*(X)), \text{ bounded, } A(g\gamma) = T_g A(\gamma) T_g^{-1}\},$$

$$\omega(A) = \int_{\Gamma_X} \int_{X/G} a_\gamma(x, x) dx \mu(d\gamma), \quad a_\gamma - \text{integral kernel of } A.$$

Theorem

\mathcal{C} is a von Neumann algebra, ω is a faithful normal semifinite trace on \mathcal{C} .

$\mathbf{P}_\gamma^{(p)}$ - ortho projection onto $\text{Ker } H_\gamma^{(p)}$. Consider maps

$\mathbf{P}^{(p)} : \gamma \mapsto \mathbf{P}_\gamma^{(p)}$, $e^{-tH^{(p)}} : \gamma \mapsto e^{-tH_\gamma^{(p)}}$. Then $\mathbf{P}^{(p)}, e^{-tH^{(p)}} \in \mathcal{C}$.

Theorem

1)

$$\omega \left(\mathbf{P}^{(p)} \right) < \infty, \quad \omega \left(e^{-tH^{(p)}} \right) < \infty$$

2) McKean-Singer formula

$$\text{STR } e^{-tH} := \sum_p (-1)^p \omega \left(e^{-tH_\gamma^{(p)}} \right) = \text{STR } \mathbf{P}$$

3)

$$\text{STR } \mathbf{P} = \chi(X/G) \text{ for any Gibbs measure } \mu$$

Proof is based on the probabilistic representation of $e^{-tH^{(p)}}$, McKean-Singer formula in general von Neumann algebras (Connes-Moscovichi) and study of short-time asymptotics of STR e^{-tH} .

Corollary

Let $\chi(X/G) \neq 0$ and μ be G -ergodic.

Then

$$\dim \text{Ker } H_\gamma = \infty \text{ for } \mu\text{-a.a. } \gamma \in \Gamma_X.$$

Open questions

- to compute individual Betti numbers
- are there measures for which the index is different?
- to consider the full Hodge-de Rham operator associated with μ on Γ_X

Recent developments- analysis of cluster measures

"Between isolated atoms or molecules and bulk materials there lies a class of unique structures, known as clusters, that consist of a few to hundreds of atoms or molecules. Within this range of "nanophase" ... materials may exhibit novel properties due to quantum confinement effects."

Dynamics of clusters: From elementary to biological structures. Po-Y. Cheng, J.S. Baskin, A.H. Zewail (Arthur Amos Noyes Laboratory of Chemical Physics, California Institute of Technology) PNAS, 2006

Recent developments- analysis of cluster measures

"Between isolated atoms or molecules and bulk materials there lies a class of unique structures, known as clusters, that consist of a few to hundreds of atoms or molecules. Within this range of "nanophase" ... materials may exhibit novel properties due to quantum confinement effects."

Dynamics of clusters: From elementary to biological structures. Po-Y. Cheng, J.S. Baskin, A.H. Zewail (Arthur Amos Noyes Laboratory of Chemical Physics, California Institute of Technology) PNAS, 2006

What is a Cluster Point Process (CPP)? (e.g. Daley, Vere-Jones)

- Take a random configuration $\gamma_c \subset X$ of *invisible* centres.
- To each $x \in \gamma_c$, attach a set $Z_x = \left\{ z_j^{(x)} \right\}_{j=1}^N$ of *observable* points (*cluster* at x).
- Resulting configuration $Z = \cup_{x \in \gamma_c} Z_x$ is a *Cluster Point Process*.

Cluster measure - distribution of CPP

Construction of cluster measures on Γ_X (L. Bogachev, AD)

① "Unpacking" mapping $p : \Gamma_{X^n} \rightarrow \Gamma_X$

$$X^n \ni \bar{x} = (x_1, \dots, x_n) \xrightarrow{p} \{x_1, \dots, x_n\} \subset X,$$

$$\Gamma_{X^n} \ni \{\dots, \bar{x}, \bar{y}, \bar{z}, \dots\} \xrightarrow{p} \{\dots, x_1, \dots, x_n, y_1, \dots, y_n, z_1, \dots, z_n, \dots\} \in \Gamma_X.$$

② Poisson or Gibbs measure μ on Γ_{X^n} .

③ Cluster measure $\mu_{cl} := p^* \mu$:

$$\mu_{cl}(A) := \mu(p(A)).$$